Schémas Bas-Froude pour le modèle Shallow Water multicouches

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Travail en collaboration avec J.P. Vila et F. Couderc

Séminaire Mathématiques et Applications - Pau

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Université Paul Sabatier



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- 2 Schéma semi-implicite
- 3 Version explicite
- 4 Résultats numériques
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6 Perspectives

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Plan

Généralités

- Présentation du modèle
- Vers un modèle régularisé
- Le modèle augementé

2 Schéma semi-implicite

- 3 Version explicite
- 4 Résultats numériques
- 5 Version mailles décalées

6 Perspectives

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Présentation du modèle

Shallow Water multicouches

$$\begin{aligned} \partial_t h_i + div(h_i \mathbf{u}_i) &= 0, \\ \partial_t(h_i \mathbf{u}_i) + div(h_i \mathbf{u}_i \otimes \mathbf{u}_i) + h_i \nabla \phi_i &= 0. \end{aligned} \qquad i = 1, \cdots, L.$$



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Présentation du modèle

Shallow Water multicouches

$$\partial_t h_i + div(h_i \mathbf{u}_i) = 0,$$

 $\partial_t (h_i \mathbf{u}_i) + div(h_i \mathbf{u}_i \otimes \mathbf{u}_i) + h_i \nabla \phi_i = 0.$

$$i=1,\cdots,L$$
 .

Difficultés :

Structure

▷ Non conservatif▷ Hyperbolicité

Cadre applicatif

 \triangleright Bas-Froude

Stabilité

- ▷ Energie
- \triangleright Robustesse
- Etats d'équilibre



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Présentation du modèle

Shallow Water multicouches

$$\partial_t h_i + div(h_i \mathbf{u}_i) = 0,$$

 $\partial_t (h_i \mathbf{u}_i) + div(h_i \mathbf{u}_i \otimes \mathbf{u}_i) + h_i \nabla \phi_i = 0.$

$$i=1,\cdots,L$$
 .

▷ Energie potentielle
$$\mathcal{E}$$
 : $\partial_{\rho_i h_i} \mathcal{E} = \phi_i = g \sum_{j=1}^{L} \frac{\rho_j h_j}{\rho_{max(i,j)}}$.
▷ Energie cinétique : $\mathcal{K}_i = \frac{1}{2} \rho_i h_i ||\mathbf{u}_i||^2$.
▷ Energie totale : $\mathcal{E} = \mathcal{E} + \sum_{i=1}^{L} \mathcal{K}_i$.

Bilans d'énergie (solutions régulières)

$$\partial_{t} \mathcal{E} + \sum_{i=1}^{L} div \left(\rho_{i} h_{i} \phi_{i} \mathbf{u}_{i} \right) - \sum_{i=1}^{L} \rho_{i} h_{i} \mathbf{u}_{i} \cdot \nabla \phi_{i} = 0,$$

$$\partial_{t} \mathcal{K}_{i} + div \left(\mathcal{K}_{i} \mathbf{u}_{i} \right) + \rho_{i} h_{i} \mathbf{u}_{i} \cdot \nabla \phi_{i} = 0,$$

$$\partial_{t} \mathcal{E} + \sum_{i=1}^{L} div \left[\left(\rho_{i} h_{i} \phi_{i} + \mathcal{K}_{i} \right) \mathbf{u}_{i} \right] = 0.$$

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Equations régularisées

Shallow Water multicouches régularisé

 $\begin{aligned} \partial_t h_i + div(h_i(\mathbf{u}_i - \delta \mathbf{u}_i)) &= 0, \\ \partial_t(h_i \mathbf{u}_i) + div(h_i \mathbf{u}_i \otimes (\mathbf{u}_i - \delta \mathbf{u}_i)) + h_i \nabla \phi_i &= 0. \end{aligned}$

$$i=1,\cdots,L$$
 .

▷ Energie potentielle
$$\mathcal{E}$$
 : $\partial_{\rho_i h_i} \mathcal{E} = \phi_i = g \sum_{j=1}^{L} \frac{\rho_j h_j}{\rho_{max}(i,j)}$
▷ Energie cinétique : $\mathcal{K}_i = \frac{1}{2} \rho_i h_i \|\mathbf{u}_i\|^2$.
▷ Energie totale : $E = \mathcal{E} + \sum_{i=1}^{L} \mathcal{K}_i$.

Bilans d'énergie (solutions régulières)

$$\partial_{t}\mathcal{E} + \sum_{i=1}^{L} div \left(\rho_{i}h_{i}\phi_{i}\mathbf{u}_{i}\right) - \sum_{i=1}^{L}\rho_{i}h_{i}\mathbf{u}_{i}.\nabla\phi_{i} = 0,$$

$$\partial_{t}\mathcal{K}_{i} + div \left(\mathcal{K}_{i}\mathbf{u}_{i}\right) + \rho_{i}h_{i}\mathbf{u}_{i}.\nabla\phi_{i} = 0,$$

$$\partial_{t}\mathcal{E} + \sum_{i=1}^{L} div \left[\left(\rho_{i}h_{i}\phi_{i} + \mathcal{K}_{i}\right)\mathbf{u}_{i}\right] = 0.$$

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Equations régularisées

Shallow Water multicouches régularisé

$$\begin{aligned} \partial_t h_i + div(h_i(\mathbf{u}_i - \delta \mathbf{u}_i)) &= 0, \\ \partial_t(h_i \mathbf{u}_i) + div(h_i \mathbf{u}_i \otimes (\mathbf{u}_i - \delta \mathbf{u}_i)) + h_i \nabla \phi_i &= 0. \end{aligned}$$

$$i=1,\cdots,L$$
 .

$$\triangleright \text{ Energie potentielle } \mathcal{E} : \partial_{\rho_i h_i} \mathcal{E} = \phi_i = g \sum_{j=1}^{L} \frac{\rho_j h_j}{\rho_{max(i,j)}}.$$

$$\triangleright \text{ Energie cinétique } : \mathcal{K}_i = \frac{1}{2} \rho_i h_i \|\mathbf{u}_i\|^2.$$

$$\triangleright \text{ Energie totale } : \mathcal{E} = \mathcal{E} + \sum_{i=1}^{L} \mathcal{K}_i.$$

Bilans d'énergie (solutions régulières)

$$\partial_{t}\mathcal{E} + \sum_{i=1}^{L} div \left(\rho_{i}h_{i}\phi_{i}(\mathbf{u}_{i}-\delta\mathbf{u}_{i})\right) - \sum_{i=1}^{L}\rho_{i}h_{i}\mathbf{u}_{i}.\nabla\phi_{i} = -\sum_{i=1}^{L}\rho_{i}h_{i}\delta\mathbf{u}_{i}.\nabla\phi_{i},$$

$$\partial_{t}\mathcal{K}_{i} + div \left(\mathcal{K}_{i}(\mathbf{u}_{i}-\delta\mathbf{u}_{i})\right) + \rho_{i}h_{i}\mathbf{u}_{i}.\nabla\phi_{i} = 0,$$

$$\partial_{t}\mathcal{E} + \sum_{i=1}^{L} div \left[\left(\rho_{i}h_{i}\phi_{i}+\mathcal{K}_{i}\right)(\mathbf{u}_{i}-\delta\mathbf{u}_{i})\right] = -\sum_{i=1}^{L}\rho_{i}h_{i}\delta\mathbf{u}_{i}.\nabla\phi_{i}.$$

Equations régularisées

Shallow Water multicouches régularisé

$$\partial_t h_i + div(h_i(\mathbf{u}_i - \delta \mathbf{u}_i)) = 0,$$

$$\partial_t(h_i\mathbf{u}_i) + div(h_i\mathbf{u}_i\otimes (\mathbf{u}_i-\delta\mathbf{u}_i)) + h_i\nabla\phi_i = 0$$
.

Bilans d'énergie (solutions régulières)

$$\partial_t \mathcal{E} + \sum_{i=1}^{L} div \left(\rho_i h_i \phi_i (\mathbf{u}_i - \delta \mathbf{u}_i) \right) - \sum_{i=1}^{L} \rho_i h_i \mathbf{u}_i . \nabla \phi_i = -\sum_{i=1}^{L} \rho_i h_i \delta \mathbf{u}_i . \nabla \phi_i ,$$

$$\partial_t \mathcal{K}_i + div \left(\mathcal{K}_i (\mathbf{u}_i - \delta \mathbf{u}_i) \right) + \rho_i h_i \mathbf{u}_i . \nabla \phi_i = 0 ,$$

$$\partial_t E + \sum_{i=1}^{L} div \left[\left(\rho_i h_i \phi_i + \mathcal{K}_i \right) (\mathbf{u}_i - \delta \mathbf{u}_i) \right] = -\sum_{i=1}^{L} \rho_i h_i \delta \mathbf{u}_i . \nabla \phi_i .$$

 $\triangleright \delta \mathbf{u}_i = \gamma \nabla \phi_i$: second membre **régularisant**.

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 $i=1,\cdots,L$

Modèle augmenté

 \triangleright On introduit la vorticité $\omega = \partial_x v - \partial_y u$ comme variable dans le système.

Modèle augmenté régularisé

$$\partial_t h_i + div(h_i(\mathbf{u}_i - \delta \mathbf{u}_i)) = 0,$$

$$\partial_t(h_i \mathbf{u}_i) + div(h_i \mathbf{u}_i \otimes (\mathbf{u}_i - \delta \mathbf{u}_i)) + h_i \nabla \phi_i = 0.$$

 \triangleright Energie totale : $E = \mathcal{E} + \sum_{i=1}^{L} \mathcal{K}_i$

Bilan d'énergie (solutions continues)

$$\partial_t E + \sum_{i=1}^L div \Big[\Big(h_i \phi_i + \mathcal{K}_i \Big) (\mathbf{u}_i - \delta \mathbf{u}_i) \Big] = - \sum_{i=1}^L \rho_i h_i \delta \mathbf{u}_i \cdot \nabla \phi_i$$

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Modèle augmenté

▷ On introduit la vorticité $\omega = \partial_x v - \partial_y u$ comme variable dans le système. ▷ On pose : $q_i = \frac{\omega_i}{h_i}$. Le système multicouches devient :

Modèle augmenté régularisé

$$\begin{aligned} \partial_t h_i + div(h_i(\mathbf{u}_i - \delta \mathbf{u}_i)) &= 0, \\ \partial_t(h_i \mathbf{u}_i) + div(h_i \mathbf{u}_i \otimes (\mathbf{u}_i - \delta \mathbf{u}_i)) + h_i \nabla \phi_i &= 0, \\ \partial_t(h_i q_i) + div(h_i q_i(\mathbf{u}_i - \delta \mathbf{u}_i)) &= 0. \end{aligned}$$

▷ Energie associée à la vorticité $W_i = \frac{1}{2}\rho_i h_i q_i^2$. ▷ Energie totale : $E_a = E + \sum_{i=1}^{L} W_i = \mathcal{E} + \sum_{i=1}^{L} \mathcal{K}_i + \sum_{i=1}^{L} W_i$.

Bilan d'énergie (solutions continues)

$$\partial_t E_a + \sum_{i=1}^{L} div \left[\left(h_i \phi_i + \mathcal{K}_i + \mathcal{W}_i \right) (\mathbf{u}_i - \delta \mathbf{u}_i) \right] = -\sum_{i=1}^{L} \rho_i h_i \delta \mathbf{u}_i \cdot \nabla \phi_i ,$$

Plan

2 Schéma semi-implicite

- Schéma numérique
- Energie potentielle
 Energie cinétique

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Schéma numérique (Parisot - Vila, 2015)

Schéma numérique

$$h_{K,i}^{n+1} = h_{K,i}^{n} - \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,K} m_{e,K}$$
$$h_{K,i}^{n+1} \mathbf{u}_{K,i}^{n+1} = h_{K,i}^{n} \mathbf{u}_{K,i}^{n} - \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \left(\mathbf{u}_{K,i}^{n} (\mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,K})^{+} - \mathbf{u}_{K_{e,i}}^{n} (\mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,K})^{-} \right) m_{e}$$
$$- \frac{\Delta t}{m_{K}} h_{K,i}^{n+1} \sum_{e \in \partial K} \Phi_{e,i}^{n+1} m_{e},$$

$$m_{\mathcal{K}} = |\mathcal{K}| \quad ; \quad m_{e} = |e|,$$
$$\omega^{\pm} = \frac{|\omega| \pm \omega}{2},$$
$$\Phi_{e,i}^{n+1} = \frac{\Phi_{\mathcal{K},i}^{n+1} + \Phi_{\mathcal{K}_{e},i}^{n+1}}{2}.$$



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Schéma numérique (Parisot - Vila, 2015)

Schéma numérique

$$h_{K,i}^{n+1} = h_{K,i}^{n} - \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,K} m_{e},$$

$$h_{K,i}^{n+1} \mathbf{u}_{K,i}^{n+1} = h_{K,i}^{n} \mathbf{u}_{K,i}^{n} - \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \left(\mathbf{u}_{K,i}^{n} (\mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,K})^{+} - \mathbf{u}_{K_{e},i}^{n} (\mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,K})^{-} \right) m_{e}$$

$$- \frac{\Delta t}{m_{K}} h_{K,i}^{n+1} \sum_{e \in \partial K} \Phi_{e,i}^{n+1} m_{e},$$

Flux numériques :

$$\mathcal{F}_{e,i}^{n+1} = h_{e,i}^{n+1} \left(\mathbf{u}_{e,i}^{n} - \gamma \Delta t \frac{\delta \Phi_{e,i}^{n+1}}{m_{e}} \right) \quad , \quad \text{où} \quad \delta \Phi_{e,i}^{n+1} = \frac{1}{2} \left(\Phi_{K_{e,i}}^{n+1} - \Phi_{K,i}^{n+1} \right) \mathbf{n}_{e,K} \,.$$

 \triangleright équivalent discret de $h_i \left(\mathbf{u}_i - \delta \mathbf{u}_i
ight)$ avec $\delta \mathbf{u}_i = \gamma
abla \phi_i$.

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Schéma sur la partie potentiel

Schéma sur \mathcal{E}_{K}^{n}

$$\mathcal{E}_{K}^{n+1} - \mathcal{E}_{K}^{n} + \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \sum_{i=1}^{L} \rho_{i} (\mathcal{G}_{\mathcal{E},e,i}^{n+1} \cdot \mathbf{n}_{e,K}) m_{e} - \mathcal{Q}_{\mathcal{E},K} \leq \mathcal{R}_{\mathcal{E},K} + \mathcal{H}_{\mathcal{E},K} .$$

$$\begin{split} \mathcal{G}_{\mathcal{E},e,i}^{n+1} \cdot \mathbf{n}_{e,K} &= \Phi_{e,i}^{n+1} \mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,K} ,\\ \mathcal{H}_{\mathcal{E},K} &= \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \sum_{i=1}^{L} \rho_{i} \left(\frac{h_{K_{e},i}^{n+1} \mathbf{u}_{K_{e},i}^{n} - h_{K,i}^{n+1} \mathbf{u}_{K,i}^{n}}{2} \right) \cdot \delta \Phi_{e,i}^{n+1} m_{e} ,\\ \mathcal{Q}_{\mathcal{E},K} &= \frac{\Delta t}{m_{K}} \sum_{i=1}^{L} \rho_{i} h_{K,i}^{n+1} \mathbf{u}_{K,i}^{n} \cdot \sum_{e \in \partial K} \Phi_{e,i}^{n+1} m_{e} ,\\ \mathcal{R}_{\mathcal{E},K} &= \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \sum_{i=1}^{L} \rho_{i} \left(\mathcal{F}_{e,i}^{n+1} - h_{e,i}^{n+1} \mathbf{u}_{e,i}^{n} \right) \cdot \delta \Phi_{e,i}^{n+1} m_{e} . \end{split}$$

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Schéma sur la partie potentiel

Schéma sur \mathcal{E}_{K}^{n}

$$\mathcal{E}_{K}^{n+1} - \mathcal{E}_{K}^{n} + \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \sum_{i=1}^{L} \rho_{i} \left(\mathcal{G}_{\mathcal{E},e,i}^{n+1} \cdot \mathbf{n}_{e,K} \right) m_{e} - \mathcal{Q}_{\mathcal{E},K} \leq \mathcal{R}_{\mathcal{E},K} + \mathcal{H}_{\mathcal{E},K} \,.$$

$$\begin{aligned} \mathcal{G}_{\mathcal{E},e,i}^{n+1} \cdot \mathbf{n}_{e,K} &= \Phi_{e,i}^{n+1} \mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,K} ,\\ \mathcal{H}_{\mathcal{E},K} &= \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \sum_{i=1}^{L} \rho_{i} \left(\frac{h_{K_{e,i}}^{n+1} \mathbf{u}_{K_{e,i}}^{n} - h_{K,i}^{n+1} \mathbf{u}_{K,i}^{n}}{2} \right) \cdot \delta \Phi_{e,i}^{n+1} m_{e} ,\\ \mathcal{Q}_{\mathcal{E},K} &= \frac{\Delta t}{m_{K}} \sum_{i=1}^{L} \rho_{i} h_{K,i}^{n+1} \mathbf{u}_{K,i}^{n} \cdot \sum_{e \in \partial K} \Phi_{e,i}^{n+1} m_{e} ,\\ \mathcal{R}_{\mathcal{E},K} &= \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \sum_{i=1}^{L} \rho_{i} \left(\mathcal{F}_{e,i}^{n+1} - h_{e,i}^{n+1} \mathbf{u}_{e,i}^{n} \right) \cdot \delta \Phi_{e,i}^{n+1} m_{e} .\end{aligned}$$

▷ Version discrète de

$$\partial_t \mathcal{E} + \sum_{i=1}^{L} div \left(\rho_i h_i \phi_i (\mathbf{u}_i - \delta \mathbf{u}_i) \right) - \sum_{i=1}^{L} \rho_i h_i \mathbf{u}_i \cdot \nabla \phi_i = -\sum_{i=1}^{L} \rho_i h_i \delta \mathbf{u}_i \cdot \nabla \phi_i.$$

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Schémas Bas-Froude SW multicouches

Schéma sur l'énergie cinétique

Schéma sur
$$\mathcal{K}_{K,i}^n = \frac{1}{2} h_{K,i}^n \|\mathbf{u}_{K,i}^n\|^2$$

$$\mathcal{K}_{K,i}^{n+1} - \mathcal{K}_{K,i}^{n} + \frac{\Delta t}{m_{\mathcal{K}}} \sum_{e \in \partial \mathcal{K}} \left(\mathcal{G}_{\mathcal{K},e,i}^{n+1} \cdot \mathbf{n}_{e,\mathcal{K}} \right) m_{e} + \mathcal{Q}_{\mathcal{K},\mathcal{K},i} \leq -\mathcal{H}_{\mathcal{K},\mathcal{K},i} + \mathcal{R}_{\mathcal{K},\mathcal{K},i} \cdot \mathbf{n}_{e,\mathcal{K},i}$$

$$\begin{aligned} \mathcal{G}_{\mathcal{K},e,i}^{n+1} \cdot \mathbf{n}_{e,\mathcal{K}} &= \frac{1}{2} \| \mathbf{u}_{\mathcal{K},i}^{n} \|^{2} \big(\mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,\mathcal{K}} \big)^{+} - \frac{1}{2} \| \mathbf{u}_{\mathcal{K}_{e},i}^{n} \|^{2} \big(\mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,\mathcal{K}} \big)^{-} , \\ \mathcal{H}_{\mathcal{K},\mathcal{K},i} &= \frac{\Delta t}{m_{\mathcal{K}}} \sum_{e \in \partial \mathcal{K}} \frac{\Delta t}{2} \left(\frac{h_{\mathcal{K}_{e},i}^{n+1}}{m_{\mathcal{K}_{e}}} - \frac{h_{\mathcal{K},i}^{n+1}}{m_{\mathcal{K}}} \right) \| \delta \Phi_{e,i}^{n+1} \|^{2} m_{e} , \\ \mathcal{Q}_{\mathcal{K},\mathcal{K},i} &= \frac{\Delta t}{m_{\mathcal{K}}} h_{\mathcal{K},i}^{n+1} \mathbf{u}_{\mathcal{K},i}^{n} \cdot \sum_{e \in \partial \mathcal{K}} \Phi_{e,i}^{n+1} m_{e} , \\ \mathcal{R}_{\mathcal{K},\mathcal{K},i} &= \frac{\Delta t}{m_{\mathcal{K}}} \sum_{e \in \partial \mathcal{K}} \Delta t \frac{h_{e,i}^{n+1}}{m_{e}^{n+1}} \| \delta \Phi_{e,i}^{n+1} \|^{2} m_{e} . \end{aligned}$$

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Schéma sur l'énergie cinétique

Schéma sur
$$\mathcal{K}_{K,i}^n = \frac{1}{2} h_{K,i}^n \|\mathbf{u}_{K,i}^n\|^2$$

$$\mathcal{K}_{K,i}^{n+1} - \mathcal{K}_{K,i}^{n} + \frac{\Delta t}{m_{\mathcal{K}}} \sum_{e \in \partial \mathcal{K}} \left(\mathcal{G}_{\mathcal{K},e,i}^{n+1} \cdot \mathbf{n}_{e,\mathcal{K}} \right) m_{e} + \mathcal{Q}_{\mathcal{K},\mathcal{K},i} \leq -\mathcal{H}_{\mathcal{K},\mathcal{K},i} + \mathcal{R}_{\mathcal{K},\mathcal{K},i} \cdot \mathbf{n}_{e,\mathcal{K},i}$$

$$\begin{aligned} \mathcal{G}_{\mathcal{K},e,i}^{n+1} \cdot \mathbf{n}_{e,\mathcal{K}} &= \frac{1}{2} \| \mathbf{u}_{\mathcal{K},i}^{n} \|^{2} \left(\mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,\mathcal{K}} \right)^{+} - \frac{1}{2} \| \mathbf{u}_{\mathcal{K}_{e,i}}^{n} \|^{2} \left(\mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,\mathcal{K}} \right)^{-}, \\ \mathcal{H}_{\mathcal{K},\mathcal{K},i} &= \frac{\Delta t}{m_{\mathcal{K}}} \sum_{e \in \partial \mathcal{K}} \frac{\Delta t}{2} \left(\frac{h_{\mathcal{K}_{e},i}^{n+1}}{m_{\mathcal{K}_{e}}} - \frac{h_{\mathcal{K},i}^{n+1}}{m_{\mathcal{K}}} \right) \| \delta \Phi_{e,i}^{n+1} \|^{2} m_{e}, \\ \mathcal{Q}_{\mathcal{K},\mathcal{K},i} &= \frac{\Delta t}{m_{\mathcal{K}}} h_{\mathcal{K},i}^{n+1} \mathbf{u}_{\mathcal{K},i}^{n} \cdot \sum_{e \in \partial \mathcal{K}} \Phi_{e,i}^{n+1} m_{e}, \\ \mathcal{R}_{\mathcal{K},\mathcal{K},i} &= \frac{\Delta t}{m_{\mathcal{K}}} \sum_{e \in \partial \mathcal{K}} \Delta t \frac{h_{e,i}^{n+1}}{m_{e}^{n+1}} \| \delta \Phi_{e,i}^{n+1} \|^{2} m_{e}. \end{aligned}$$

 $\triangleright \text{ Version discrète de } \partial_t \mathcal{K}_i + div \left(\mathcal{K}_i (\mathbf{u}_i - \delta \mathbf{u}_i) \right) + h_i \mathbf{u}_j, \nabla \phi_j = \mathbf{0}.$

Energie totale

$$E^{n+1} - E^n = \sum_{K} \sum_{i=1}^{L} m_K \rho_i (\mathcal{K}_{K,i}^{n+1} - \mathcal{K}_{K,i}^n) + \sum_{K} m_K (\mathcal{E}_K^{n+1} - \mathcal{E}_K^n)$$

$$\leq \sum_{K} \left(\sum_{i=1}^{L} m_K \rho_i \mathcal{R}_{\mathcal{K},K,i} + m_K \mathcal{R}_{\mathcal{E},K} \right)$$

$$\leq \Delta t \sum_{K} \sum_{i=1}^{L} \sum_{e \in \partial K} \rho_i (1 - \gamma) \Delta t \frac{h_{e,i}^{n+1}}{m_e^{n+1}} \| \delta \Phi_{e,i}^{n+1} \|^2 m_e.$$

 \triangleright Décroissance de l'énergie numérique $\longrightarrow \gamma \geq 1$.

▷ Résultat identique avec le modèle augmenté (pas d'impact sur la production d'énergie).

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Plan

Généralités

2 Schéma semi-implicite

3 Version explicite

- Le schéma
- Définition des termes régularisants
- Analyse asymptotique

4) Résultats numériques

5 Version mailles décalées

6 Perspectives

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Le schéma

Schéma numérique

$$h_{K,i}^{n+1} = h_{K,i}^{n} - \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \mathcal{F}_{e,i}^{n} \cdot \mathbf{n}_{e,K} m_{e,K}$$

$$h_{K,i}^{n+1} \mathbf{u}_{K,i}^{n+1} = h_{K,i}^{n} \mathbf{u}_{K,i}^{n} - \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \left(\mathbf{u}_{K,i}^{n} (\mathcal{F}_{e,i}^{n} \cdot \mathbf{n}_{e,K})^{+} - \mathbf{u}_{K_{e,i}}^{n} (\mathcal{F}_{e,i}^{n} \cdot \mathbf{n}_{e,K})^{-} \right) m_{e}$$

$$- \frac{\Delta t}{m_{K}} h_{K,i}^{n} \sum_{e \in \partial K} \Phi_{e,i}^{n,*} \mathbf{n}_{e,K} m_{e} \cdot \mathbf{n}_{e,K}$$

$$\begin{split} m_{K} &= |K| \quad ; \quad m_{e} = |e| , \\ \omega^{\pm} &= \frac{1}{2} \left(|\omega| \pm \omega \right) , \\ \Phi_{e,i}^{n} &= \frac{1}{2} \left(\Phi_{K,i}^{n} + \Phi_{K_{e},i}^{n} \right) . \end{split}$$



Le schéma

Schéma numérique

$$h_{K,i}^{n+1} = h_{K,i}^{n} - \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \mathcal{F}_{e,i}^{n} \cdot \mathbf{n}_{e,K} m_{e},$$

$$h_{K,i}^{n+1} \mathbf{u}_{K,i}^{n+1} = h_{K,i}^{n} \mathbf{u}_{K,i}^{n} - \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \left(\mathbf{u}_{K,i}^{n} (\mathcal{F}_{e,i}^{n} \cdot \mathbf{n}_{e,K})^{+} - \mathbf{u}_{K_{e},i}^{n} (\mathcal{F}_{e,i}^{n} \cdot \mathbf{n}_{e,K})^{-} \right) m_{e}$$

$$- \frac{\Delta t}{m_{K}} h_{K,i}^{n} \sum_{e \in \partial K} \Phi_{e,i}^{n,*} \mathbf{n}_{e,K} m_{e}.$$

$$\begin{split} m_{K} &= |K| \quad ; \quad m_{e} = |e| ,\\ \omega^{\pm} &= \frac{1}{2} \left(|\omega| \pm \omega \right) ,\\ \Phi_{e,i}^{n} &= \frac{1}{2} \left(\Phi_{K,i}^{n} + \Phi_{K_{e},i}^{n} \right) . \end{split}$$

• $\mathcal{F}_{e,i}^n=(h\mathbf{u})_{e,i}^n-\Pi_{e,i}^n$.



Le schéma

Schéma numérique

$$h_{K,i}^{n+1} = h_{K,i}^{n} - \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \mathcal{F}_{e,i}^{n} \cdot \mathbf{n}_{e,K} m_{e},$$

$$h_{K,i}^{n+1} \mathbf{u}_{K,i}^{n+1} = h_{K,i}^{n} \mathbf{u}_{K,i}^{n} - \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \left(\mathbf{u}_{K,i}^{n} (\mathcal{F}_{e,i}^{n} \cdot \mathbf{n}_{e,K})^{+} - \mathbf{u}_{K_{e},i}^{n} (\mathcal{F}_{e,i}^{n} \cdot \mathbf{n}_{e,K})^{-} \right) m_{e}$$

$$- \frac{\Delta t}{m_{K}} h_{K,i}^{n} \sum_{e \in \partial K} \Phi_{e,i}^{n,*} \mathbf{n}_{e,K} m_{e}.$$

 $e \in \partial K \cap \partial K_e$

 M_K

K

 $\mathbf{n}_{e,K}$

$$\begin{split} m_{K} &= |K| \quad ; \quad m_{e} = |e| , \\ \omega^{\pm} &= \frac{1}{2} \left(|\omega| \pm \omega \right) , \\ \Phi_{e,i}^{n} &= \frac{1}{2} \left(\Phi_{K,i}^{n} + \Phi_{K_{e},i}^{n} \right) . \end{split}$$

• $\mathcal{F}_{e,i}^{n} = (h\mathbf{u})_{e,i}^{n} - \prod_{e,i}^{n}$. • $\Phi_{e,i}^{n,*} = \Phi_{e,i}^{n} - \Lambda_{e,i}^{n}$. M_{K_e}

Termes régularisants

Flux régularisés

$$\mathcal{F}_{e,i}^{n} = (h\mathbf{u})_{e,i}^{n} - \Pi_{e,i}^{n} = (h\mathbf{u})_{e,i}^{n} - \gamma\Delta t(h\mu)_{e,i}^{n+1/2}\delta\Phi_{e,i}^{n} \to \text{correction en } \nabla\phi$$
$$\text{avec } (h\mu)_{e,i}^{n+1/2} = \frac{1}{2} \left(\frac{(h_{K,i}^{n})^{2}}{h_{K,i}^{n+1}} \frac{m_{\partial K}}{m_{K}} + \frac{(h_{K_{e},i}^{n})^{2}}{h_{K_{e},i}^{n+1}} \frac{m_{\partial K_{e}}}{m_{K_{e}}} \right).$$

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Termes régularisants

Flux régularisés

$$\mathcal{F}_{e,i}^{n} = (h\mathbf{u})_{e,i}^{n} - \Pi_{e,i}^{n} = (h\mathbf{u})_{e,i}^{n} - \gamma\Delta t(h\mu)_{e,i}^{n+1/2}\delta\Phi_{e,i}^{n} \to \text{correction en } \nabla\phi$$
$$\text{avec } (h\mu)_{e,i}^{n+1/2} = \frac{1}{2} \left(\frac{(h_{K,i}^{n})^{2}}{h_{K,i}^{n+1}} \frac{m_{\partial K}}{m_{K}} + \frac{(h_{Ke,i}^{n})^{2}}{h_{Ke,i}^{n+1}} \frac{m_{\partial Ke}}{m_{Ke}} \right).$$

Potentiel régularisé

$$\begin{split} \Phi_{e,i}^{n,*} = & \Phi_{e,i}^n - \Lambda_{e,i}^n = \Phi_{e,i}^n - \alpha \nu \Delta t \mu_e \delta(h\mathbf{u})_{e,i}^n \quad \to \text{ correction en } \nabla.h\mathbf{u} \\ \text{avec } \mu_e = & \frac{1}{2} \left(\frac{m_{\partial K}}{m_K} + \frac{m_{\partial K_e}}{m_{K_e}} \right) \,. \end{split}$$

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Termes régularisants

Flux régularisés

$$\mathcal{F}_{e,i}^{n} = (h\mathbf{u})_{e,i}^{n} - \Pi_{e,i}^{n} = (h\mathbf{u})_{e,i}^{n} - \gamma\Delta t(h\mu)_{e,i}^{n+1/2} \delta\Phi_{e,i}^{n} \to \text{correction en } \nabla\phi$$
$$\text{avec } (h\mu)_{e,i}^{n+1/2} = \frac{1}{2} \left(\frac{(h_{K,i}^{n})^{2}}{h_{K,i}^{n+1}} \frac{m_{\partial K}}{m_{K}} + \frac{(h_{K_{e,i}}^{n})^{2}}{h_{K_{e,i}}^{n+1}} \frac{m_{\partial K_{e}}}{m_{K_{e}}} \right).$$

Potentiel régularisé

$$\begin{split} \Phi_{e,i}^{n,*} = & \Phi_{e,i}^n - \Lambda_{e,i}^n = \Phi_{e,i}^n - \alpha \nu \Delta t \mu_e \delta(h\mathbf{u})_{e,i}^n \quad \to \text{correction en } \nabla.h\mathbf{u} \\ \text{avec } \mu_e = & \frac{1}{2} \left(\frac{m_{\partial K}}{m_K} + \frac{m_{\partial K_e}}{m_{K_e}} \right) \,. \end{split}$$

Conditions sur γ, α

$$\gamma$$
 , $lpha\in\mathcal{S}:=\left\{x$, $p(x)=(k_e)^2x^2-x+1\leq 0
ight\}$

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Cas 1d barotrope

Conditions

$$\gamma, \alpha \in \mathcal{S} := \left\{ x, p(x) = (k_e)^2 x^2 - x + 1 \le 0 \right\} \quad , \quad k_e = 2 \frac{\Delta t}{\Delta x} c_e \, .$$

$$\begin{split} \Delta &= 1 - 4k_e^2 \ge 0 \Rightarrow \frac{\Delta t}{\Delta x} c_e \le \frac{1}{4} \text{ (Condition CFL)}, \\ x^{\pm} &= \frac{1 \pm \sqrt{1 - 4k_e^2}}{2k_e^2}, \\ \mathcal{S} &= \begin{bmatrix} x^-, x^+ \end{bmatrix}, \text{ condition à priori locale.} \end{split}$$

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Cas 1d barotrope

Conditions

$$\gamma, \alpha \in \mathcal{S} := \left\{ x, p(x) = (k_e)^2 x^2 - x + 1 \le 0 \right\} \quad , \quad k_e = 2 \frac{\Delta t}{\Delta x} c_e \, .$$

$$\begin{split} \Delta &= 1 - 4k_e^2 \ge 0 \Rightarrow \frac{\Delta t}{\Delta x} c_e \le \frac{1}{4} \text{ (Condition CFL)} \\ x^{\pm} &= \frac{1 \pm \sqrt{1 - 4k_e^2}}{2k_e^2}, \\ \mathcal{S} &= \begin{bmatrix} x^-, x^+ \end{bmatrix}, \text{ condition à priori locale.} \end{split}$$

Choix des paramètres de diffusion

 $\label{eq:alpha} \begin{array}{l} \triangleright \ 2 \in \mathcal{S} : \alpha = \gamma = 2 \ \text{permet d'assurer la stabilité du schéma.} \\ \triangleright \ \text{Expérimentation numérique : schémas trop diffusifs en pratique .} \\ \rightarrow \ \text{On cherche des conditions relaxées en régime bas-Froude :} \end{array}$

$$\alpha = \gamma = 1/2$$

 \rightarrow Analyse de stabilité linéaire : importance capitale du schéma temporel.

Cartographies de stabilité

Schémas d'ordre 1 et RK2 en temps



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Cartographies de stabilité

Numérique Vs Analyse linéaire. Ordre 1 et MUSCL en espace



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Interprétation semi discrète

$$\begin{cases} h_i^{n+1} - h_i^n = -\Delta t \partial_x \left(hu_i\right)^n + \left(\Delta t\right)^2 \gamma \partial_x \left(h_i \frac{\partial_x \phi_i}{\epsilon^2}\right)^n \\ hu_i^{n+1} - hu_i^n = -\Delta t \left(\partial_x \left(\bar{u}_i hu_i^*\right)\right)^n \\ -\Delta t \left(h_i \frac{\partial_x \phi_i}{\epsilon^2}\right)^n + \left(\Delta t\right)^2 \alpha \left(h_i \frac{\partial_{xx} \left(hu_i\right)}{\epsilon^2}\right)^r \end{cases}$$

avec $hu_i^* = hu_i - \Delta t \gamma \left(h_i \frac{\partial_x \phi_i}{\epsilon^2}\right)$ et $\bar{u}_i = u_i + \mathcal{O}(\Delta x)$.

Analyse temps court
$$t = \epsilon \tau$$

Equation de la masse

Interprétation semi discrète

$$\begin{cases} h_i^{n+1} - h_i^n = -\Delta t \partial_x (hu_i)^n + (\Delta t)^2 \gamma \partial_x \left(h_i \frac{\partial_x \phi_i}{\epsilon^2} \right)^n \\ hu_i^{n+1} - hu_i^n = -\Delta t \left(\partial_x \left(\bar{u}_i hu_i^* \right) \right)^n \\ -\Delta t \left(h_i \frac{\partial_x \phi_i}{\epsilon^2} \right)^n + (\Delta t)^2 \alpha \left(h_i \frac{\partial_{xx} (hu_i)}{\epsilon^2} \right)^n \end{cases}$$

avec $hu_i^* = hu_i - \Delta t \gamma \left(h_i \frac{\partial_x \phi_i}{\epsilon^2} \right)$ et $\bar{u}_i = u_i + \mathcal{O}(\Delta x)$.

Analyse temps court
$$t = \epsilon \tau$$

Equation de la masse

$$\frac{h_{i}^{n+1}-2h_{i}^{n}+h_{i}^{n-1}}{\left(\Delta\tau\right)^{2}}=-\frac{\epsilon}{\Delta\tau}\left[\partial_{x}\left(hu_{i}^{n}-hu_{i}^{n-1}\right)\right]+\gamma\left[\partial_{x}\left(\left(h_{i}\partial_{x}\phi_{i}\right)^{n}-\left(h_{i}\partial_{x}\phi_{i}\right)^{n-1}\right)\right]$$

Interprétation semi discrète

$$\begin{cases} h_i^{n+1} - h_i^n = -\Delta t \partial_x \left(hu_i\right)^n + \left(\Delta t\right)^2 \gamma \partial_x \left(h_i \frac{\partial_x \phi_i}{\epsilon^2}\right)^n \\ hu_i^{n+1} - hu_i^n = -\Delta t \left(\partial_x \left(\bar{u}_i hu_i^*\right)\right)^n \\ -\Delta t \left(h_i \frac{\partial_x \phi_i}{\epsilon^2}\right)^n + \left(\Delta t\right)^2 \alpha \left(h_i \frac{\partial_{xx} \left(hu_i\right)}{\epsilon^2}\right)^r \end{cases}$$

avec $hu_i^* = hu_i - \Delta t \gamma \left(h_i \frac{\partial_x \phi_i}{\epsilon^2}\right)$ et $\bar{u}_i = u_i + \mathcal{O}(\Delta x)$.

Analyse temps court
$$t = \epsilon \tau$$

Equation de la masse

$$\frac{h_{i}^{n+1}-2h_{i}^{n}+h_{i}^{n-1}}{\left(\Delta\tau\right)^{2}}=-\frac{\epsilon}{\Delta\tau}\left[\partial_{x}\left(hu_{i}^{n}-hu_{i}^{n-1}\right)\right]+\underbrace{\gamma\left[\partial_{x}\left(\left(h_{i}\partial_{x}\phi_{i}\right)^{n}-\left(h_{i}\partial_{x}\phi_{i}\right)^{n-1}\right)\right]}_{\mathcal{O}(\Delta\tau)}$$

Interprétation semi discrète

$$\begin{cases} h_i^{n+1} - h_i^n = -\Delta t \partial_x (hu_i)^n + (\Delta t)^2 \gamma \partial_x \left(h_i \frac{\partial_x \phi_i}{\epsilon^2} \right)^n \\ hu_i^{n+1} - hu_i^n = -\Delta t \left(\partial_x \left(\bar{u}_i hu_i^* \right) \right)^n \\ -\Delta t \left(h_i \frac{\partial_x \phi_i}{\epsilon^2} \right)^n + (\Delta t)^2 \alpha \left(h_i \frac{\partial_{xx} (hu_i)}{\epsilon^2} \right)^n \end{cases}$$

avec
$$hu_i^* = hu_i - \Delta t \gamma \left(h_i \frac{\partial_x \varphi_i}{\epsilon^2} \right)$$
 et $\bar{u}_i = u_i + \mathcal{O}(\Delta x)$.

Analyse temps court $t = \epsilon \tau$

$$\frac{h_i^{n+1} - 2h_i^n + h_i^{n-1}}{(\Delta \tau)^2} = -\frac{\epsilon}{\Delta \tau} \left[\partial_x \left(h u_i^n - h u_i^{n-1} \right) \right] + \mathcal{O}(\Delta \tau).$$

$$\frac{\epsilon}{\Delta \tau} \left(h u_i^n - h u_i^{n-1} \right) = -\epsilon^2 \left(\partial_x \left(\bar{u}_i h u_i^* \right) \right)^{n-1}$$

$$-\epsilon^2 \left(h_i \frac{\partial_x \phi_i}{\epsilon^2} \right)^{n-1} + \epsilon \Delta \tau \alpha \left(h_i \partial_{xx} \left(h u_i \right) \right)^{n-1}$$

Interprétation semi discrète

$$\begin{cases} h_i^{n+1} - h_i^n = -\Delta t \partial_x \left(hu_i\right)^n + \left(\Delta t\right)^2 \gamma \partial_x \left(h_i \frac{\partial_x \phi_i}{\epsilon^2}\right)^n \\ hu_i^{n+1} - hu_i^n = -\Delta t \left(\partial_x \left(\bar{u}_i hu_i^*\right)\right)^n \\ -\Delta t \left(h_i \frac{\partial_x \phi_i}{\epsilon^2}\right)^n + \left(\Delta t\right)^2 \alpha \left(h_i \frac{\partial_{xx} \left(hu_i\right)}{\epsilon^2}\right)^n \end{cases}$$

avec
$$hu_i^* = hu_i - \Delta t \gamma \left(h_i \frac{\partial_x \phi_i}{\epsilon^2} \right)$$
 et $\bar{u}_i = u_i + \mathcal{O}(\Delta x)$.

Analyse temps court $t = \epsilon \tau$

$$\frac{h_i^{n+1} - 2h_i^n + h_i^{n-1}}{(\Delta \tau)^2} = -\frac{\epsilon}{\Delta \tau} \left[\partial_x \left(h u_i^n - h u_i^{n-1} \right) \right] + \mathcal{O}(\Delta \tau).$$
$$\frac{\epsilon}{\Delta \tau} \left(h u_i^n - h u_i^{n-1} \right) = -\epsilon^2 \left(\partial_x \left(\bar{u}_i h u_i^* \right) \right)^{n-1} - \left(h_i \partial_x \phi_i \right)^n + \mathcal{O}(\Delta \tau)$$

Interprétation semi discrète

$$\begin{cases} h_i^{n+1} - h_i^n = -\Delta t \partial_x \left(hu_i\right)^n + \left(\Delta t\right)^2 \gamma \partial_x \left(h_i \frac{\partial_x \phi_i}{\epsilon^2}\right)^n \\ hu_i^{n+1} - hu_i^n = -\Delta t \left(\partial_x \left(\bar{u}_i hu_i^*\right)\right)^n \\ -\Delta t \left(h_i \frac{\partial_x \phi_i}{\epsilon^2}\right)^n + \left(\Delta t\right)^2 \alpha \left(h_i \frac{\partial_{xx} \left(hu_i\right)}{\epsilon^2}\right)^n \end{cases}$$

avec
$$hu_i^* = hu_i - \Delta t \gamma \left(h_i \frac{\partial_x \varphi_i}{\epsilon^2} \right)$$
 et $\bar{u}_i = u_i + \mathcal{O}(\Delta x)$.

Analyse temps court $t = \epsilon \tau$

$$\frac{h_i^{n+1} - 2h_i^n + h_i^{n-1}}{(\Delta \tau)^2} = -\frac{\epsilon}{\Delta \tau} \left[\partial_x \left(h u_i^n - h u_i^{n-1} \right) \right] + \mathcal{O}(\Delta \tau).$$
$$\frac{\epsilon}{\Delta \tau} \left(h u_i^n - h u_i^{n-1} \right) = -\epsilon^2 \left(\partial_x \left(\bar{u}_i h u_i^* \right) \right)^{n-1} - \left(h_i \partial_x \phi_i \right)^n + \mathcal{O}(\Delta \tau).$$

$$\left. \begin{array}{l} \epsilon^{2}hu_{i}^{*} = \mathcal{O}_{\epsilon \to 0}(\epsilon^{2}) + \mathcal{O}(\Delta \tau) \\ \bar{u}_{i} = \mathcal{O}(1) \end{array} \right\} \Rightarrow \epsilon^{2} \left(\partial_{x} \left(\bar{u}_{i}hu_{i}^{*} \right) \right)^{n-1} = \begin{array}{c} \mathcal{O}_{\epsilon \to 0}(\epsilon^{2}) + \mathcal{O}(\Delta \tau) \\ \epsilon \to 0 \end{array}$$

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Interprétation semi discrète

$$\begin{cases} h_i^{n+1} - h_i^n = -\Delta t \partial_x (hu_i)^n + (\Delta t)^2 \gamma \partial_x \left(h_i \frac{\partial_x \phi_i}{\epsilon^2} \right)^n \\ hu_i^{n+1} - hu_i^n = -\Delta t \left(\partial_x \left(\bar{u}_i h u_i^* \right) \right)^n \\ -\Delta t \left(h_i \frac{\partial_x \phi_i}{\epsilon^2} \right)^n + (\Delta t)^2 \alpha \left(h_i \frac{\partial_{xx} (hu_i)}{\epsilon^2} \right)^n \end{cases}$$
$$hu_i^* = hu_i - \Delta t \gamma \left(h_i \frac{\partial_x \phi_i}{\epsilon^2} \right) \text{ et } \bar{u}_i = u_i + \mathcal{O}(\Delta x) .$$

Analyse temps court $t = \epsilon \tau$

$$\frac{h_i^{n+1}-2h_i^n+h_i^{n-1}}{\left(\Delta\tau\right)^2}=\partial_x\left(h_i\partial_x\phi_i\right)^n+\underset{\varepsilon\to 0}{\mathcal{O}}(\epsilon^2)+\mathcal{O}(\Delta\tau)\,.$$

avec

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Plan

Généralités

2 Schéma semi-implicite

3 Version explicite

Résultats numériques

- Ondes de gravité 2d en régime linéaire
- Transport de vortex sur β plan

Version mailles décalées

6 Perspectives

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Ondes de gravité 2d en régime linéaire

Configuration

- 5 couches .
- $\rho_i = 1000 + 10(i-1)$.
- $\ell = 100$ km, h0 = 5000 m.

Condition initiale

Perturbation de surface : • $\delta h(\cos(kx) + \cos(ky))$,

$$\delta h = 1m, k = \frac{2\pi}{\ell}.$$

Paramètres numériques

- grille 41×41 .
- $\alpha = \beta = 0.1$.
- Ordre 2 temps (RK2) et espace (MUSCL) .



Transport de vortex sur β plan : calibration

Paramètres

L = 2	$R_{earth} = 6367442,76m$
$h_0 = 5000 m$	$\Omega = 2\pi/(3600/24)$
$h_1 = 2500m$	$f_0=2\Omega sin(heta)\simeq$ 9, 25410 $^{-5}$
$ heta=38.5^\circ$	$\beta = 2\Omega cos(\theta) / R_{earth} \simeq 1,78810^{-11}$

- Domaine [-1000, 1000] × [-1000, 1000]
- Condition initiale déterminée à partir de la distribution de densité.

Coriolis : splitting

Approximation β plan : $f = f_0 + \beta y$.

$$u^{n+1} = \cos(f\Delta t)u^n + \sin(f\Delta t)v^n,$$

$$v^{n+1} = \cos(f\Delta t)v^n - \sin(f\Delta t)u^n.$$

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Transport de vortex sur β plan : déviation sur 100 jours



(a) $\Delta x = 2km$



(d) $\Delta x = 20 km$

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(c) $\Delta x = 10 km$

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Schémas Bas-Froude SW multicouches

Transport de vortex sur β plan : déviation sur 100 jours



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Plan

- 1 Généralités
- 2 Schéma semi-implicite
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 Contexte
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 - Bilan d'énergie

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Extension aux mailles décalées

Formalisme : R. Herbin, W. Kheriji, J.C. Latché, K.Saleh (2013, 2014).

Schéma

$$\begin{split} h_{K,i}^{n+1} &= h_{K,i}^{n} - \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,K} m_{e} \\ h_{D^{e},i}^{n+1} \mathbf{u}_{D^{e},i}^{n+1} &= h_{D^{e},i}^{n} \mathbf{u}_{D^{e},i}^{n} - \frac{\Delta t}{m_{D^{e}}} \sum_{f \in \partial D^{e}} \left(\mathbf{u}_{D^{e},i}^{n} \left(\mathcal{F}_{f,i}^{n+1} \cdot \mathbf{n}_{f,D^{e}} \right)^{+} - \mathbf{u}_{D_{f}}^{n} \left(\mathcal{F}_{f,i}^{n+1} \cdot \mathbf{n}_{f,D^{e}} \right)^{-} \right) m_{f} \\ &- \frac{\Delta t}{m_{D^{e}}} h_{D^{e},i}^{n+1} \nabla \Phi_{e,i} \, . \end{split}$$



Extension aux mailles décalées

Formalisme : R. Herbin, W. Kheriji, J.C. Latché, K.Saleh (2013, 2014).

Schéma

$$h_{K,i}^{n+1} = h_{K,i}^{n} - \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,K} m_{e}$$

$$h_{D^{e},i}^{n+1} \mathbf{u}_{D^{e},i}^{n+1} = h_{D^{e},i}^{n} \mathbf{u}_{D^{e},i}^{n} - \frac{\Delta t}{m_{D^{e}}} \sum_{f \in \partial D^{e}} \left(\mathbf{u}_{D^{e},i}^{n} (\mathcal{F}_{f,i}^{n+1} \cdot \mathbf{n}_{f,D^{e}})^{+} - \mathbf{u}_{D_{f}^{e},i}^{n} (\mathcal{F}_{f,i}^{n+1} \cdot \mathbf{n}_{f,D^{e}})^{-} \right) m_{f}$$

$$- \frac{\Delta t}{m_{D^{e}}} h_{D^{e},i}^{n+1} \nabla \Phi_{e,i} \cdot$$

▷ Flux numériques : $\mathcal{F}_{e,i}^{n+1} = h_{D^e,i}^{n+1} (\mathbf{u}_{D^e,i}^n - \delta \mathbf{u}_{D^e,i}^n)$. ▷ Gradient de pression : $\Phi_{e,i}$ à déterminer.



Extension aux mailles décalées

Formalisme : R. Herbin, W. Kheriji, J.C. Latché, K.Saleh (2013, 2014).

Schéma

$$h_{K,i}^{n+1} = h_{K,i}^{n} - \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,K} m_{e}$$

$$h_{D^{e},i}^{n+1} \mathbf{u}_{D^{e},i}^{n+1} = h_{D^{e},i}^{n} \mathbf{u}_{D^{e},i}^{n} - \frac{\Delta t}{m_{D^{e}}} \sum_{f \in \partial D^{e}} \left(\mathbf{u}_{D^{e},i}^{n} (\mathcal{F}_{f,i}^{n+1} \cdot \mathbf{n}_{f,D^{e}})^{+} - \mathbf{u}_{D_{f}^{e},i}^{n} (\mathcal{F}_{f,i}^{n+1} \cdot \mathbf{n}_{f,D^{e}})^{-} \right) m_{f}$$

$$- \frac{\Delta t}{m_{D^{e}}} h_{D^{e},i}^{n+1} \nabla \Phi_{e,i} \cdot$$

 \triangleright Flux numériques : $\mathcal{F}_{e_i}^{n+1} = h_{D^e_i}^{n+1} (\mathbf{u}_{D^e_i}^n - \delta \mathbf{u}_{D^e_i}^n)$. \triangleright Gradient de pression : $\Phi_{e,i}$ à déterminer.

Hauteur d'eau auxiliaire $|D^{e}|h_{D^{e}i}^{n+1} = |D_{K}^{e}|h_{Ki}^{n+1} + |D_{K_{e}i}^{e}|h_{K_{ei}i}^{n+1},$ $D_K^e \boldsymbol{\ell}$ $h_{D^{\mathbf{e}},i}^{n+1} - h_{D^{\mathbf{e}},i}^{n} = -\frac{\Delta t}{m_{D^{\mathbf{e}}}} \sum_{f \in \Delta D^{\mathbf{e}}} \mathcal{F}_{f,i}^{n+1} \cdot \mathbf{n}_{f,D^{\mathbf{e}}} m_{f}.$ $f \in \partial D^e$

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Schémas d'énergie : partie cinétique

Schéma sur
$$\mathcal{K}_{D^e,i}^n = \frac{1}{2} h_{D^e,i}^n \|\mathbf{u}_{D^e,i}^n\|^2$$

$$\mathcal{K}_{D^e,i}^{n+1} - \mathcal{K}_{D^e,i}^n \leq -\frac{\Delta t}{m_{D^e}} \sum_{f \in \partial D^e} \left(\mathcal{G}_{\mathcal{K},f,i}^{n+1} \cdot \mathbf{n}_{f,D^e} \right) m_f - \mathcal{Q}_{\mathcal{K},e,i} + \mathcal{R}_{\mathcal{K},e,i} \, .$$

$$\begin{aligned} \mathcal{G}_{\mathcal{K},f,i}^{n+1} \cdot \mathbf{n}_{f,D^{e}} &= \frac{1}{2} \| \mathbf{u}_{D^{e},i}^{n} \|^{2} (\mathcal{F}_{f,i}^{n+1} \cdot \mathbf{n}_{f,D^{e}})^{+} - \frac{1}{2} \| \mathbf{u}_{D_{f}^{e},i}^{n} \|^{2} (\mathcal{F}_{f,i}^{n+1} \cdot \mathbf{n}_{f,D^{e}})^{-}, \\ \mathcal{Q}_{\mathcal{K},e,i} &= \Delta t \, h_{D^{e},i}^{n+1} \mathbf{u}_{D^{e},i}^{n} (\nabla \Phi)_{e,i}, \\ \mathcal{R}_{\mathcal{K},e,i} &= h_{D^{e},i}^{n+1} (\Delta t)^{2} \| (\nabla \Phi)_{e,i} \|^{2}. \end{aligned}$$

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Schémas d'énergie : partie potentiel

Schéma sur ${\cal E}$

$$\mathcal{E}_{K}^{n+1} - \mathcal{E}_{K}^{n} \leq -\frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \sum_{i=1}^{L} \left(\mathcal{G}_{\mathcal{E},e,i}^{n} \cdot \mathbf{n}_{e,K} \right) m_{e} + \mathcal{Q}_{\mathcal{E},K} - \mathcal{R}_{\mathcal{E},K} \,.$$

$$\mathcal{G}_{\mathcal{E},e,i}^{n} \cdot \mathbf{n}_{e,K} = \Phi_{e,i}^{n+1} \mathcal{F}_{e,i}^{n+1} \cdot \mathbf{n}_{e,K} ,$$

$$\mathcal{Q}_{\mathcal{E},K} = \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \sum_{i=1}^{L} h_{D^{e},i}^{n+1} \mathbf{u}_{e,i}^{n} \cdot \delta \Phi_{e,i}^{n+1} m_{e} ,$$

$$\mathcal{R}_{\mathcal{E},K} = \frac{\Delta t}{m_{K}} \sum_{e \in \partial K} \sum_{i=1}^{L} h_{D^{e},i}^{n+1} \delta \mathbf{u}_{e,i}^{n} \cdot \delta \Phi_{e,i}^{n+1} m_{e}$$

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Energie totale $E^{n+1} - E^n \leq \sum_{D^e} \sum_{i=1}^{L} \rho_i m_{D^e} (\mathcal{K}_{D^e,i}^{n+1} - \mathcal{K}_{D^e,i}^n) + \sum_{K} m_K (\mathcal{E}_K^{n+1} - \mathcal{E}_K^n).$

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Energie totale $E^{n+1} - E^n \leq \sum_{D^e} \sum_{i=1}^{L} \rho_i m_{D^e} \left(-Q_{\mathcal{K}, e, i} + \mathcal{R}_{\mathcal{K}, e, i} \right) \\ + \sum_{\mathcal{K}} m_{\mathcal{K}} \left(Q_{\mathcal{E}, \mathcal{K}} - \mathcal{R}_{\mathcal{E}, \mathcal{K}} \right) .$

Energie totale

$$E^{n+1} - E^n \leq \sum_{D^e} \sum_{i=1}^{L} \rho_i m_{D^e} \left(-\mathcal{Q}_{\mathcal{K},e,i} + \mathcal{R}_{\mathcal{K},e,i} \right) \\ + \sum_{\mathcal{K}} m_{\mathcal{K}} \left(\mathcal{Q}_{\mathcal{E},\mathcal{K}} - \mathcal{R}_{\mathcal{E},\mathcal{K}} \right) .$$

 \triangleright Termes en Q :

$$\sum_{D^{e}} \sum_{i=1}^{L} \rho_{i} m_{D^{e}} \mathcal{Q}_{\mathcal{K},e,i} = \Delta t \sum_{D^{e}} \sum_{i=1}^{L} \rho_{i} m_{D^{e}} h_{D^{e},i}^{n+1} (\nabla \Phi)_{e,i} \mathbf{u}_{D^{e},i}^{n}$$
$$\sum_{\mathcal{K}} m_{\mathcal{K}} \mathcal{Q}_{\mathcal{E},\mathcal{K}} = \Delta t \sum_{\mathcal{K}} \sum_{e \in \partial \mathcal{K}} \sum_{i=1}^{L} \rho_{i} h_{D^{e},i}^{n+1} \mathbf{u}_{e,i}^{n} \cdot \delta \Phi_{e,i}^{n+1} m_{e}$$

Energie totale

$$E^{n+1} - E^n \leq \sum_{D^e} \sum_{i=1}^{L} \rho_i m_{D^e} \left(-\mathcal{Q}_{\mathcal{K},e,i} + \mathcal{R}_{\mathcal{K},e,i} \right) \\ + \sum_{\mathcal{K}} m_{\mathcal{K}} \left(\mathcal{Q}_{\mathcal{E},\mathcal{K}} - \mathcal{R}_{\mathcal{E},\mathcal{K}} \right) .$$

 \triangleright Termes en Q :

$$\sum_{D^{e}} \sum_{i=1}^{L} \rho_{i} m_{D^{e}} \mathcal{Q}_{\mathcal{K},e,i} = \Delta t \sum_{D^{e}} \sum_{i=1}^{L} \rho_{i} m_{D^{e}} h_{D^{e},i}^{n+1} (\nabla \Phi)_{e,i} \mathbf{u}_{D^{e},i}^{n}$$
$$\sum_{K} m_{K} \mathcal{Q}_{\mathcal{E},K} = \Delta t \sum_{K} \sum_{e \in \partial K} \sum_{i=1}^{L} \rho_{i} h_{D^{e},i}^{n+1} \mathbf{u}_{e,i}^{n} . \delta \Phi_{e,i}^{n+1} m_{e}$$

Compensation des termes : $(\nabla \Phi)_{e,i} = \frac{m_e}{m_{D^e}} \left(\Phi_{K_e,i}^{n+1} - \Phi_{K,i}^{n+1} \right) \mathbf{n}_{e,K}.$

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Energie totale

$$E^{n+1} - E^n \leq \sum_{D^e} \sum_{i=1}^{L} \rho_i m_{D^e} \left(-Q_{\mathcal{K}, e, i} + \mathcal{R}_{\mathcal{K}, e, i} \right) \\ + \sum_{\mathcal{K}} m_{\mathcal{K}} \left(Q_{\mathcal{E}, \mathcal{K}} - \mathcal{R}_{\mathcal{E}, \mathcal{K}} \right) .$$

 \triangleright Termes en \mathcal{R} :

$$\sum_{D^{e}} \sum_{i=1}^{L} m_{D^{e}} \mathcal{R}_{\mathcal{K},e,i} = \sum_{D^{e}} \sum_{i=1}^{L} m_{D^{e}} h_{D^{e},i}^{n+1} (\Delta t)^{2} \| (\nabla \Phi)_{e,i} \|^{2}$$
$$\sum_{\mathcal{K}} m_{\mathcal{K}} \mathcal{R}_{\mathcal{E},\mathcal{K}} = \Delta t \sum_{\mathcal{K}} \sum_{e \in \partial \mathcal{K}} \sum_{i=1}^{L} h_{D^{e},i}^{n+1} \delta \mathbf{u}_{e,i}^{n} . \delta \Phi_{e,i}^{n+1} m_{e}$$

Energie totale

$$E^{n+1} - E^n \leq \sum_{D^e} \sum_{i=1}^{L} \rho_i m_{D^e} \left(-\mathcal{Q}_{\mathcal{K}, e, i} + \mathcal{R}_{\mathcal{K}, e, i} \right) \\ + \sum_{\mathcal{K}} m_{\mathcal{K}} \left(\mathcal{Q}_{\mathcal{E}, \mathcal{K}} - \mathcal{R}_{\mathcal{E}, \mathcal{K}} \right) .$$

 \triangleright Termes en \mathcal{R} :

$$\sum_{D^{e}} \sum_{i=1}^{L} m_{D^{e}} \mathcal{R}_{\mathcal{K},e,i} = \sum_{D^{e}} \sum_{i=1}^{L} m_{D^{e}} h_{D^{e},i}^{n+1} (\Delta t)^{2} \| (\nabla \Phi)_{e,i} \|^{2}$$
$$\sum_{\mathcal{K}} m_{\mathcal{K}} \mathcal{R}_{\mathcal{E},\mathcal{K}} = \Delta t \sum_{\mathcal{K}} \sum_{e \in \partial \mathcal{K}} \sum_{i=1}^{L} h_{D^{e},i}^{n+1} \delta \mathbf{u}_{e,i}^{n} \cdot \delta \Phi_{e,i}^{n+1} m_{e}$$

Restes négatifs : $\delta \mathbf{u}_{e,i}^n = 2\gamma \frac{m_e}{m_{D^e}} \Delta t \delta \Phi_{e,i}^{n+1}$, $\gamma \ge 1$.

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- Généralités
- 2 Schéma semi-implicite
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- 5 Version mailles décalées
- 6 Perspectives

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Perspectives

Schéma explicite

• Stabilité.

 \rightarrow Convergence des analyses (linéaire / non linéaire, semi-discret / discret / continu).

- \rightarrow Ordre élevé temps et espace Ordre très élevé (MOOD).
- \rightarrow Influence de Coriolis.
- Version augmentée.

Perspectives

Schéma explicite

- Stabilité.
- Version augmentée.

Mailles décalées

- Implémentation.
- Consistance.
 - \rightarrow Cas MAC et Rannacher Turek (1D , 2D).
 - \rightarrow Approches DF, et variationnelle.
- Analyse asymptotique.

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Perspectives

Schéma explicite

- Stabilité.
- Version augmentée.

Mailles décalées

- Implémentation.
- Consistance.
- Analyse asymptotique.

Conditions aux limites transparentes

R. Monjarret - The multi-layer shallow water model with free surface. Numerical treatment of the open boundaries (thèse - 2014).

- Validation sur le schéma explicite.
- Expérimentation dans le cas implicite.
- Implémentation sur mailles décalées.

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Perspectives

Schéma explicite

- Stabilité.
- Version augmentée.

Mailles décalées

- Implémentation.
- Consistance.
- Analyse asymptotique.

Conditions aux limites transparentes

- Validation sur le schéma explicite.
- Expérimentation dans le cas implicite.
- Implémentation sur mailles décalées.

Autres perspectives opérationnelles

- Schéma hybride explicite/implicite : découplage barocline / barotrope.
- Disparition des couches, fronts secs.

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Merci !

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