# Convergence to equilibrium for semigroups in Hilbert space 

An intertwining approach

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## Part I. <br> Abstract Framework

## Contraction and Markov semigroups

Let $\mathcal{H}$ be a Hilbert space.

- We say $P=\left(P_{t}\right)_{t \geqslant 0}$ is a contraction semigroup if it is a $C_{0}$-semigroup on $\mathcal{H}$ and satisfies $\left\|\mid P_{t}\right\| \| \leqslant 1$ for all $t \geqslant 0$.


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- We say $P$ is an $L^{2}(\mu)$-Markov semigroup if it is a Markov semigroup with invariant probability measure $\mu$, i.e.

$$
\mu P_{t} f=\mu f:=\int f d \mu
$$

## Convergence to equilibrium

Let $P_{\infty}$ be the projection onto $\left\{f \in \mathcal{H}: P_{t} f=f\right.$ for all $\left.t \geqslant 0\right\}$. Our aim is to understand the long-time behavior of

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Note: If $\mathbf{P}$ is another projection such that $\left\|P_{t} f-\mathbf{P} f\right\| \rightarrow 0$, then by invariance $\mathbf{P}=P_{\infty}$.

## Spectral gap inequality

$P$ satisfies a spectral gap inequality $S(\mathfrak{a})$ if

$$
\left\|P_{t} f-P_{\infty} f\right\|_{\mathcal{H}} \leqslant e^{-\mathfrak{a} t}\left\|f-P_{\infty} f\right\|_{\mathcal{H}}
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where $\mathfrak{a}>0$ is the spectral gap of $P$ (i.e. the smallest such constant).

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- Bakry, Gentil, Ledoux's monograph [BGL14]
- Miclo [Mic15] self-adjoint, ergodic and hyperbounded $\Longrightarrow$ spectral gap


## Hypocoercivity \& perturbed spectral gap

$P$ satisfies a hypocoercive estimate $\mathrm{H}(C, \omega)$ if

$$
\left\|P_{t} f-P_{\infty} f\right\|_{\mathcal{H}} \leqslant C e^{-\omega t}\left\|f-P_{\infty} f\right\|_{\mathcal{H}},
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- Villani's memoir [Vil09] (Lyapunov functional techniques)
- Mischler, Mouhot [MM16] (shrinking \& enlarging spaces)
- Baudoin [Bau13] (Bakry-Émery theory \& 「-calculus)
- Hérau \& Nier [HN04] (hypoellipticity techniques)


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$P$ satisfies a perturbed spectral gap inequality $\operatorname{PS}(C, \omega)$ if it satisfies $\mathrm{H}(C, \omega)$ and $\omega$ is a gap in the point spectrum.
- Patie \& Savov [PS16] (intertwining relationships).


## Intertwining

## Definition 2

Let $P$ and $Q$ be contraction semigroups on Hilbert spaces $\mathcal{H}$ and $\mathcal{K}$.

- We say that $P$ intertwines with $Q$, written $P \triangleleft Q$, if there exists $\Lambda \in \mathbf{B}(\mathcal{K}, \mathcal{H})$ such that on $\mathcal{K}$ and for all $t \geqslant 0$ we have

$$
P_{t} \Lambda=\Lambda Q_{t}
$$

- If $P \triangleleft Q$ and $\Lambda$ is bijective then $P_{t}=\Lambda Q_{t} \Lambda^{-1}$. In this case we say $P$ is in the similarity orbit of $Q$, and write $P \bowtie Q$.


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Literature:

- Sz.-Nagy \& Foais [SNFBK10] (quasi-affinity)
- Douglas [Dou69] (unitary equivalence for normal operators)
- Dynkin [Dyn82] (Dynkin's criterion)
- Rogers \& Pitman [RP80] (Brownian motion \& Bessel process)


## First convergence result

Theorem 1 (Patie, V. 2016)
Suppose $P \bowtie Q$ and let $\kappa(\Lambda)=\| \| \Lambda\| \|\left\|\Lambda^{-1}\right\| \geqslant 1$.
(i) If $Q$ satisfies $H(C, \omega)$ then $P$ satisfies $H(C \kappa(\Lambda), \omega)$.
(ii) If $Q$ satisfies $S(\mathfrak{a})$ then $P$ satisfies $P S(\kappa(\Lambda), \mathfrak{a})$.

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Proof.
Let $P \bowtie Q$. Then, $P_{\infty}=\Lambda Q_{\infty} \Lambda^{-1}$. Since $Q$ satisfies $\mathrm{H}(C, \omega)$,

$$
\begin{aligned}
\left\|P_{t} f-P_{\infty} f\right\|_{\mathcal{H}} & =\left\|\Lambda Q_{t} \Lambda^{-1} f-\Lambda Q_{\infty} \Lambda^{-1} f\right\|_{\mathcal{H}} \\
& \leqslant\|\Lambda \Lambda\|\| \| Q_{t} \Lambda^{-1} f-Q_{\infty} \Lambda^{-1} f \|_{\mathcal{K}} \\
& \leqslant C\| \| \Lambda\left\|e^{-\omega t}\right\| \Lambda^{-1} f-Q_{\infty} \Lambda^{-1} f \|_{\mathcal{K}} \\
& =C\| \| \Lambda\left\|e^{-\omega t}\right\| \Lambda^{-1} f-\Lambda^{-1} P_{\infty} f \|_{\mathcal{K}} \\
& \leqslant C\|\Lambda \Lambda\|\| \| \Lambda^{-1}\| \| e^{-\omega t}\left\|f-P_{\infty} f\right\|_{\mathcal{H}} \\
& =C \kappa(\Lambda) e^{-\omega t}\left\|f-P_{\infty} f\right\|_{\mathcal{H}}
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## Beyond similarity via resolution of identity

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By the spectral theorem we have

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Q_{t}=\int_{\sigma\left(-A_{Q}\right)} e^{-\gamma t} d E_{\gamma},
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where $E$ is the resolution of the identity for $-A_{Q}$, the generator of $Q$.

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Writing $g=f-Q_{\infty} f$, we have

$$
\begin{aligned}
\left\|Q_{t} g\right\|^{2} & =\int_{\sigma\left(-A_{Q}\right) \backslash\{0\}} e^{-2 \Re(\gamma) t} d\left\langle E_{\gamma} g, g\right\rangle \\
& \leqslant e^{-2 \mathfrak{a} t} \int_{\sigma\left(-A_{Q}\right) \backslash\{0\}} d\left\langle E_{\gamma} g, g\right\rangle \\
& \leqslant e^{-\mathfrak{a} t}\|g\|^{2} .
\end{aligned}
$$

## Non-self-adjoint (unbounded) resolutions

From now on, suppose $P \triangleleft Q$ and the intertwining operator $\Lambda$ satisfies
(a) $\operatorname{Ran}(\Lambda) \subset_{d} \mathcal{H},(\overline{\operatorname{Ran}(\Lambda)}=\mathcal{H})$
(b) $\operatorname{Ker}\left(\Lambda \upharpoonright_{D}\right)=0$ where $\bigcup_{B \in \mathcal{B}(\mathbb{C})} \operatorname{Ran}\left(E_{B}\right) \subset \mathcal{D} \subset_{d} \mathcal{H}$,

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Then on $\mathcal{D}$ we can define, for $B \in \mathcal{B}(\mathbb{C})$,

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F_{B}=\Lambda E_{B} \Lambda^{-1}
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Each $F_{B}$ is unbounded and non-self-adjoint.

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## Proposition 1

Under the above assumptions, for each $f \in \mathcal{D}$ and all $t \geqslant 0$,

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As a corollary we deduce that for each $f \in \mathcal{D}$ and all $t \geqslant 0$,

$$
\left\|P_{t} f-P_{\infty} f\right\|_{\mathcal{H}} \leqslant\|\Lambda\|\left\|e^{-\mathfrak{a} t}\right\| \Lambda^{-1}\left(f-P_{\infty} f\right) \|_{\mathcal{K}} .
$$

Cannot be extended by density because $\Lambda^{-1}$ is not necessarily bounded.

## Integral representation of $P$

Assume further that there exists $m: \sigma\left(-A_{Q}\right) \rightarrow \mathbb{C}$ such that
(c) $F_{\infty}^{m}=\int_{\sigma\left(-A_{Q}\right)} m(\gamma) d F_{\gamma}$ is bounded on $\mathcal{H}$,
(d) for $t>T \geqslant 0, u(t)=\sup _{\Re(\gamma) \geqslant a} \frac{e^{-\Re(\gamma) t}}{|m(\gamma)|}<+\infty$.

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Theorem 2 (Patie, V. 2016)
Suppose $\Lambda$ satisfies (a)-(d). Then for all $f \in \mathcal{H}$ and $t>T$,
(i) $P_{t} f=\int_{\sigma\left(-A_{Q}\right)} \frac{e^{-\gamma t}}{m(\gamma)} d F_{\gamma}^{m} f$,
(ii) $\left\|P_{t} f-P_{\infty} f\right\|_{\mathcal{H}} \leqslant\left\|F_{\infty}^{m}\right\|\|(t)\| f-P_{\infty} f \|_{\mathcal{H}}$,

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(ii) $\left\|P_{t} f-P_{\infty} f\right\|_{\mathcal{H}} \leqslant\left\|F_{\infty}^{m}\right\| u(t)\left\|f-P_{\infty} f\right\|_{\mathcal{H}}$,

## Remark 1

- If the sup is attained at $\mathfrak{a}$ then we have hypocoercivity estimate with spectral explanation for the constant, which is the spectral gap of $Q$.
- $F_{\infty}^{m}$ is no longer a resolution of the identity.
- Some sufficient conditions can be obtained when $P \triangleleft Q$ and $Q \triangleleft P$.


## Proof of Theorem 2

The proof of (i) uses Proposition 1, condition (c) and the density of $\mathcal{D}$.
For the proof of (ii), writing $g=f-P_{\infty} f$, we have

$$
\begin{aligned}
\left\|P_{t} g\right\|^{2} & =\int_{\sigma\left(-A_{Q}\right) \backslash\{0\}} \frac{e^{-\gamma t}}{m(\gamma)} d\left\langle F_{\gamma}^{m} g, P_{t} g\right\rangle \\
& \leqslant u(t) \int_{\sigma\left(-A_{Q}\right) \backslash\{0\}} d\left\langle F_{\gamma}^{m} g, P_{t} g\right\rangle \\
& \leqslant\left\|F_{\infty}^{m}\right\| u(t)\|g\|\left\|P_{t} g\right\|
\end{aligned}
$$

## Part II. <br> Generalized Laguerre semigroups

## Characterization

Let $\mathcal{N}$ denote the set of all functions $\psi: i \mathbb{R} \rightarrow \mathbb{C}$ given by

$$
\Psi(z)=\sigma^{2} z^{2}+m z+\int_{\mathbb{R}}\left(e^{z y}-1-z y\right) \Pi(d y)
$$

where $\sigma^{2}, m \geqslant 0$ and $\Pi$ is a measure satisfying $\int_{\mathbb{R}}\left(|x| \wedge x^{2}\right) \Pi(d x)<\infty$.

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$P=\left(e^{-t A}\right)_{t \geqslant 0} \in \mathcal{G} \mathcal{L}$, the set of generalized Laguerre semigroups, if

$$
\mathcal{M}_{A f}(z+1)=-\Psi(-z) \mathcal{M}_{f}(z)+(z+1) \mathcal{M}_{f}(z+1), \quad z \in i \mathbb{R},
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with $\mathcal{M}_{f}(s)=\int_{0}^{\infty} x^{s-1} f(x) d x$ the Mellin transform of $f$.

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with $\mathcal{M}_{f}(s)=\int_{0}^{\infty} x^{s-1} f(x) d x$ the Mellin transform of $f$.
There is a bijection between $\mathcal{N}$ and $\mathcal{G L}$. One can show that for suitable functions $f$,

$$
\begin{aligned}
-A f(x)= & \sigma^{2} x f^{\prime \prime}(x)+\left(m+\sigma^{2}-x\right) f^{\prime}(x) \\
& +\int_{\mathbb{R}}\left(f\left(e^{y} x\right)-f(x)-y x f^{\prime}(x)\right) \frac{\Pi(d y)}{x} .
\end{aligned}
$$

## Properties and classical Laguerre

Every $P \in \mathcal{G \mathcal { L }}$ is an $\mathrm{L}^{2}(\mu)$-Markov semigroup with state space $\left(\mathbb{R}_{+}, \mathcal{B}\right)$ and absolutely continuous, invariant probability measure $\mu$.

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$P$ is self-adjoint on $L^{2}(\mu)$ if and only if $\Pi \equiv 0$.
If $\sigma=1, \Pi \equiv 0$ we get the classical Laguerre semigroup of order $m \geqslant 0$, denoted by $Q^{(m)}$, which admits a spectral expansion

$$
Q_{t}^{(m)} f=\sum_{n=0}^{\infty} e^{-n t}\left\langle f, e_{n}^{(m)}\right\rangle e_{n}^{(m)}
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From this we deduce that,

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\left\|Q_{t}^{(m)} f-Q_{\infty}^{(m)} f\right\|_{L^{2}(\mu)} \leqslant e^{-t}\left\|f-Q_{\infty}^{(m)} f\right\|_{L^{2}(\mu)} .
$$

## The similarity case

Theorem 3 (Patie, V. 2016)
Let $P \in \mathcal{G} \mathcal{L}$ and suppose that the associated $\psi \in \mathcal{N}$ satisfies

$$
\Psi(z)=\Psi(-z) \Longleftrightarrow m=0 \quad \text { and } \quad \Pi(d x) \mathbb{1}_{\{x<0\}}=\Pi(d x) \mathbb{1}_{\{x>0\}}
$$

Then $P \bowtie Q^{(0)}$, where the intertwining operator $\wedge$ is a Mellin convolution operator with explicit Mellin multiplier. Furthermore,

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\left\|P_{t} f-P_{\infty} f\right\|_{L^{2}(\mu)} \leqslant \kappa(\Lambda) e^{-t}\left\|f-P_{\infty} f\right\|_{L^{2}(\mu)} .
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## Remark 2

- When $\Pi \not \equiv 0$ then $P$ is non-local and non-self-adjoint.
- The generator may or may not have a diffusive component.


## Hypocoercivity for generalized Laguerre

Wiener-Hopf factorization: $\Psi(z)=-\phi_{-}(z) \phi_{+}(-z), z \in i \mathbb{R}$, where

$$
\begin{aligned}
& \phi_{ \pm}(z)=\mathrm{k}_{ \pm}+\mathrm{d}_{ \pm} z+\int_{0}^{\infty}\left(1-e^{-z y}\right) \eta_{ \pm}(d y)=\mathrm{k}_{ \pm}+\mathrm{d}_{ \pm} z+\int_{0}^{\infty} e^{-z y} \bar{\eta}_{ \pm}(y) d y \\
& \mathrm{k}_{ \pm}, \mathrm{d}_{ \pm} \geqslant 0 \text { and } \int_{0}^{\infty}(1 \wedge y) \eta_{ \pm}(d y)<\infty
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\end{aligned}
$$

Theorem 4 (Patie, V. 2016)
Let $P \in \mathcal{G \mathcal { L }}$. If $\Psi$ is such that $\sigma>0$ and $\bar{\eta}_{-} \stackrel{0}{\sim} \bar{\eta}_{+}$, then

$$
\left\|P_{t} f-P_{\infty} f\right\|_{L^{2}(\mu)} \leqslant C e^{-t}\left\|f-P_{\infty} f\right\|_{L^{2}(\mu)}
$$

holds for all $f \in L^{2}(\mu)$ and $t \geqslant 0$, where $C>1$ is explicit.

## Sketch of proof

Key idea: $P \triangleleft Q^{(a)}$ and $Q^{(b)} \triangleleft P$ with $0 \leqslant a<b$, i.e.

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The two intertwinings allow us to conclude that

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F_{\gamma}=\sum_{n \leqslant \gamma}\left\langle\Lambda^{-1} f, e_{n}^{(a)}\right\rangle \Lambda e_{n}^{(a)}=\sum_{n \leqslant \gamma}\left\langle f, V_{n}\right\rangle P_{n},
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where $V_{n}, P_{n} \in \mathrm{~L}^{2}(\mu)$.

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$$

where $V_{n}, P_{n} \in \mathrm{~L}^{2}(\mu)$. The function $m$ is given by

$$
m(n)=\frac{\Gamma(b+1)}{\Gamma(a+1)} \frac{\Gamma(n+a+1)}{\Gamma(n+b+1)},
$$

which behaves asymptotically like $n^{a-b}$.

## Conclusions

- Intertwining seems to be a fruitful approach for investigating convergence to equilibrium.
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- Useful for transferring (spectral) information from known to unknown objects.


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How far can one go with intertwining?

## Thank you for your time and attention

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